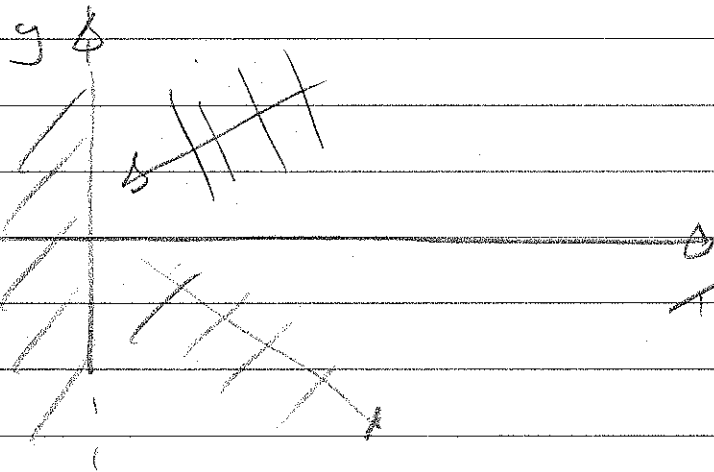


# A Short Course in "Boundary Condition Theory". PART I



$$\begin{cases}
 U_t = AU_x + BU_y & ; \quad U = \begin{pmatrix} \psi \\ \alpha_x \\ \alpha_y \end{pmatrix} \quad \begin{matrix} U \rightarrow 0 \\ y \rightarrow \pm\infty \end{matrix} \\
 \mathcal{L}U = g \\
 U(x, y, 0) = f(x, y)
 \end{cases}$$

$x \geq 0, \quad -\infty < y < \infty$   
 $+ \geq 0$

$$\left( A = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad B = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \right)$$

## Energy-method

$$\int_{-\infty}^{\infty} \int_0^{\infty} U U_t \, dx \, dy = \int_{-\infty}^{\infty} \int_0^{\infty} (U A U)_x \, dx \, dy + \int_{-\infty}^{\infty} \int_0^{\infty} \frac{1}{2} (U B U)_y \, dx \, dy$$

$$\frac{1}{2} \frac{d}{dt} \|U\|_t^2 = - \int_{-\infty}^{\infty} U^T A U \, dy$$

= 0

2)

How many boundary conditions?

Eigenvalues to A?

$$\begin{vmatrix} -\lambda & 1 & 0 \\ 1 & -\lambda & 0 \\ 0 & 0 & -\lambda \end{vmatrix} = -\lambda(\lambda^2 - 1) = 0 \Rightarrow$$

$$\lambda = 0, \lambda = 1, \lambda = -1$$

∴ One negative  $\Rightarrow$  1 B.C. = 1 eq.

$LU = q$  has 1 row.

Note: Not 1 x 3 b.c's = 3 eq's

↑  
components

What conditions?

$$-\int_{-\infty}^{\infty} u^T A u \, dy = -\int_{-\infty}^{\infty} z u \, dx \, dy$$

choose for example:

$$v = \alpha x + y \Rightarrow L = \begin{bmatrix} 1 & -\alpha & 0 \end{bmatrix}$$

3)

$$\therefore \frac{d}{dt} \|u\|^2 = - \int_{-\infty}^{\infty} 2\alpha x (\alpha x + y) dy$$

Choose by putting  $q = \alpha x + y = 0$ .

if  $\alpha < 0$  : Well-posed!

Strongly well-posed?  $q \neq 0$

$$\begin{aligned} \text{Let } \alpha > 0 &\Rightarrow 2\alpha x + 2\alpha x q = \\ &= \frac{1}{2\alpha} q^2 + 2\alpha x + 2\alpha x q - \frac{1}{2\alpha} q^2 = \\ &= -\frac{1}{2\alpha} q^2 + (\sqrt{2\alpha} x + \sqrt{2\alpha} q)^2 \Rightarrow \end{aligned}$$

$$\begin{aligned} \frac{d}{dt} \|u\|^2 &= - \int_{-\infty}^{\infty} \frac{q^2}{2\alpha} + (\sqrt{2\alpha} x + \sqrt{2\alpha} q)^2 dy \leq \\ &= - \frac{1}{2\alpha} \int_{-\infty}^{\infty} q^2 dy \end{aligned}$$

$\therefore$  Strongly well-posed.  
for  $\alpha > 0$ .

4)

## Normal Mode Analysis

$$u = e^{st + iux + iwy} \Rightarrow (u = \text{periodic})$$

$$(SI - A - i\omega B)\psi = 0, \text{ non-triv.}$$

$$\text{solution } \psi \neq 0 \Rightarrow$$

$$\det(SI - A - i\omega B) = 0 \Rightarrow$$

$$S(u + s\omega) = 0.$$

Special case  $s=0$ , steady case.

$$\Rightarrow (u_x)_x + (u_y)_y = 0, \text{ div}(u_x, u_y) = 0$$

$$u_x = 0$$

$$\Rightarrow u = \text{const.}$$

$$u_y = 0$$

$u = 0$  ( $x=0$ ) seems a good b.c.

$$5) \quad S \neq 0 \quad \Rightarrow$$

$$k_{\pm} = \pm \sqrt{S^2 + \omega^2} \quad \begin{array}{l} k_-^2 < 0 \Rightarrow 1 \text{ b.c. at } x=0 \\ k_+^2 > 0 \Rightarrow 1 \text{ b.c. at } x=1 \end{array}$$

$$\underline{u = u_+} \Rightarrow \begin{pmatrix} S - k_+ & -i\omega \\ -k_+ & S \\ -i\omega & 0 & S \end{pmatrix} \psi_+ = 0$$

$$\Rightarrow \psi_+ = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \frac{k_+}{S} \\ \frac{i\omega}{S} \end{pmatrix}$$

$$u = u_- \Rightarrow \psi_- = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \frac{k_-}{S} \\ \frac{i\omega}{S} \end{pmatrix}$$

$$\vec{u} = d_+ \psi_+ e^{k_+ x} + d_- \psi_- e^{k_- x}$$

only 2 unknowns

$\Rightarrow$   
only 2 b.c.s.

Note only 2 totally and  
if 1 boundary only one

6)

We consider one boundary at

$$x=0 \Rightarrow d_x = 0$$

$$\vec{u} = d_x \psi_x e^{i\omega x}$$

$$v - \alpha d_x^2 = \vec{g}^T \Rightarrow$$

$$d_x \psi_x^{(1)} - \alpha d_x \psi_x^{(2)} = \vec{g}^T \Rightarrow$$

$$\underbrace{(\psi_x^{(1)} - \alpha \psi_x^{(2)})}_{C(s)} d_x = \vec{g}^T$$

$$C(s) = \left(1 - \alpha \frac{u_x}{s}\right) = \left(1 + \alpha \sqrt{s^2 + \omega^2}\right) =$$

$$1 + \alpha \sqrt{1 + \left(\frac{|\omega|}{s}\right)^2}, \text{ Note } s \neq 0.$$

$$\frac{|\omega|}{s} = \frac{|\omega|(y - i\frac{1}{2})}{y^2 + \frac{1}{4}} = \frac{|\omega|y - i\frac{1}{2}|\omega|}{y^2 + \frac{1}{4}} = \frac{y}{y^2 + \frac{1}{4}} - i\frac{1}{2} \frac{|\omega|}{y^2 + \frac{1}{4}}$$

$$\therefore C(s) = 1 + \alpha \sqrt{1 + s^2}$$

7)

obviously for  $\text{Real } s > 0$  also  
 $\text{Real } (\xi) > 0$ .

$C(\xi) \neq 0$  for  $\text{Real } (\xi) > 0$

$\therefore$  Laplace's ok. no exp. growth.

How about Kreiss cond.?

Let  $\xi = i\eta \Rightarrow$

$$C = 1 + \alpha \sqrt{1 - \eta^2}$$

1)  $\frac{\eta^2}{4} < 1 \Rightarrow C > 0$

2)  $\frac{\eta^2}{4} = 1 \Rightarrow C = 1 \neq 0$

3)  $\frac{\eta^2}{4} > 1 \Rightarrow C = 1 + \alpha i \sqrt{\eta^2 - 4} \neq 0$

No possibility for failure

$$\therefore |d| \leq |\eta| / |C|$$

8) Note also  $\alpha=0$  ok

since in that case

$$C = 1 \Rightarrow d = g^2 \Rightarrow$$

$$u = g e^{kx} \Rightarrow$$

$$\vec{u}(x, \omega, s) \leq \vec{q}(\omega, s) + \text{Parabola}$$

$$\int_0^T |u|^2 dt \leq \int_0^T q^2 dt$$

Very Nice result.



9)

A short Course in Weak  
Implementation of Boundary  
Conditions, Part II

$$U_1 = AU_1 + Bf$$

$$LU = g \quad x \geq 0$$

$$U(x, y, 0) = f(x)$$

$$d = [1 - \alpha \ 0] \quad U = \begin{pmatrix} U \\ d_x \\ d_y \end{pmatrix}$$

We skip the  $y$ -dependence for now.

$$U_1 = (\tilde{P}^T \otimes A) U + \tilde{P}^T (e_0 \times \tilde{\Gamma}_0) (LU - g)$$

$$\tilde{U}(x) = f$$

position in space

from PDE

Energy  $\Rightarrow$

to be determined

$$U^T (\tilde{P} \otimes A) U_1 = U^T (C \otimes A) U + U_0^T \tilde{\Gamma}_0^T (LU - g)$$

size of  $A = 3 \times 3$

10)

Add transpose  $\Rightarrow$ 

$$\left( \begin{matrix} u \\ u \\ u \end{matrix} \right)_{\text{pos}}^T + = u^T (C \alpha C^T \oplus A) + u_0 (P_0 \alpha + \alpha^T P_0) u$$

collect terms at  $i=0$

$$- u_0^T P_0 \eta - \eta^T P_0^T u_0$$

ignore from the start.

$$B^T = u_0^T (-A + P_0 \alpha + \alpha^T P_0) u_0 =$$

 $\hat{A}$   $L = \text{is a matrix.}$ 

$$\begin{pmatrix} T_1 \\ T_2 \\ T_3 \end{pmatrix} (1 - \alpha \quad 0) = \begin{pmatrix} T_1 - \alpha T_1 & 0 \\ T_2 - \alpha T_2 & 0 \\ T_3 - \alpha T_3 & 0 \end{pmatrix}$$

+ Transpose  $\Rightarrow$ 

$$\hat{A} = \begin{pmatrix} 0 & -1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} -\alpha T_1 & T_2 - \alpha T_1 & T_3 \\ T_2 - \alpha T_1 & -\alpha T_2 & -\alpha T_3 \\ T_3 & -\alpha T_3 & 0 \end{pmatrix}$$

$$\hat{L} = 0 \Rightarrow T_1 = T_3 = 0 \quad T_2 = 1 \Rightarrow$$

$$\hat{A} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -\alpha & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

n/

How about data? Use it

So finally we get

$$\begin{aligned} BT &= -2\alpha d_\alpha^2 - \alpha_0^T \alpha_0 g - g^T \alpha_0 \alpha_0 = -2\alpha d_\alpha^2 - \alpha_0 g \\ &= \frac{g^2}{2\alpha} - \frac{g^2}{2\alpha} - 2\alpha d_\alpha^2 - \alpha_0 g = \\ & \frac{g^2}{2\alpha} (\text{rad} + \text{rad} g) \end{aligned}$$

∴ exactly as in PDF case.