Lecture 4

1 Initial Boundary Value Problem "The Big Picture"



B.C. ? Where, how many, on what form?

$$IBVP \begin{cases} u_t + P(u, (\partial/\partial x))u &= F(x, t), \quad x \in \Omega \\ Lu &= g(x, t), \quad x \in \delta\Omega \\ u(x, 0) &= f(x), \quad x \in \Omega \end{cases}$$

IVP = No Boundaries, Periodic. No second row above.

IBVP "Roughly Speaking"

$$P + (L) \longrightarrow \tilde{P} \qquad U_t + \tilde{P}u = \tilde{F}$$

$$\Rightarrow \qquad (IVP)$$

$$F + (g) \longrightarrow \tilde{F} \qquad u(x,0) = f$$

 $\tilde{P},\,\tilde{F}$ generalized operator, data. Eigenvalue analysis (shown later) $\Rightarrow\,\tilde{P}u=(\Lambda^R+i\Lambda^I)u$

Hyperbolic	$\Lambda^R\approx 0$	(transport, Euler, Maxwell, b.c.=?)
Parabolic	$\Lambda^R > 0$	(damping, heat, diffusion, b.c. everywhere)
Incompletely Parabolic	$\Lambda^R \geq 0$	(N-S, mixed systems, b.c.=?)
Well-posed	$ \Lambda^R \leq const.$	b.c. !!

- i) Must choose L such that $P + L = \tilde{P}$, OK and do not cause explosion. P often given and OK.
- ii) Need to choose L such that we have data Lu-g=0.
- iii) i) and ii) often in conflict.

 $\begin{array}{l} \underline{\mathbf{EX}}:\\ u=u_{\infty}\\ u_{x}=0\\ \alpha u+\beta u_{x}=u_{\infty} \end{array}$



2 Initial Value Problem + Fourier Expansion. (IVP)

• Continuous + semi discrete \Rightarrow

$$(\hat{u}_w)_t = \hat{P}(iw)\hat{u}_w + \hat{F} \qquad (*)$$
$$\hat{u}_w(0) = \hat{f}$$

$$\hat{u}_w(t) = e^{\hat{P}(iw)t}\hat{f} + e^{\hat{P}(iw)t}\int_0^t e^{-\hat{P}(iw)\zeta}\hat{F}d\zeta \qquad (**)$$

Knowledge about $\hat{P}(iw)$ +Parsevals inequality \Rightarrow Well-posedness or Stability

Easy to analyze since:

• $\frac{\hat{P}(iw)$ "Small" matrix The Fourier Modes "decouple the problem".

• Eigenvalues, eigenvector possible to compute analytically.

3 Initial Boundary Value Problem (IBVP)

- Continuous: <u>cannot be</u> formulated as an ODE, see (*) and (**) above.
- Semi-discrete: <u>can be</u> put up as a semi-discrete system.

$$u_t = Au + F$$

$$(* * *)$$

$$u(0) = f$$

However, the matrix A (corresponding to $\hat{P}(iw)$) is not small.

• <u>A is not "small"</u>. No "decoupling", all grid points included.

• Eigenvalues, eigenvectors almost impossible to compute analytically. Also b.c. often "ruin" the structure of A.

• Other methods and techniques are necessary.

4 Different Levels of Approximations

- i) Complicated non-linear system of eq's with initial and b.c.'s Ex: N-S with Shocks \approx Impossible to analyze
- ii) Simplified problem (model problem) with similar character. Possible to analyze
- iii) Iterate between i) and ii)

Model Problems



Maxwell, Euler, Multi D, Non-linear

Wave propogation, easier



5 Initial Boundary Value Problems. Semidiscrete Approximation "The method of lines"

 $\underline{\mathbf{EX}}$:

$$\frac{d}{dt}u_{j} = D_{0}u_{j}, \quad j = 1, ..., N - 1$$

$$u_{0} = 2u_{1} - u_{2}$$

$$u_{N} = g(t)$$

$$u_{j}(0) = f_{j}, \quad j = 0, ..., N$$
(2.28)

Nh=1, Linear extrapolation, as numerical boundary condition, Not physical!!

Various forms of (2.28), possible variants include (2.29)-(2.31), see the book.

$$\left\{ \begin{array}{ccc} \frac{du_1}{dt} = \frac{u_2 - (2u_1 - u_2)}{2\Delta x} = \frac{u_2 - u_1}{\Delta x} & & \frac{d}{dt}u = Qu + F \\ & \Rightarrow & \\ \frac{du_{N-1}}{dt} = \frac{g - (u_{N-2})}{2h} = -\frac{u_{N-2}}{\Delta x} + \frac{g}{2h} & & u(0) = f \end{array} \right\}$$
(2.19)

The general form is:

$$\frac{du_j}{dt} = Qu_j + F_j, \quad j = 1, ..., N - 1
B_n u = g
u_j(0) = f_j, \qquad j = 0, ..., N$$
(2.32)

j = 1, ..., N-1, inner points. $B_n u = g$. Complete set of boundary conditions (real + numerical).

 $B_n u = g$ include as many conditions that are needed for the ODE system to have a unique solution. The <u>number of boundary conditions</u> is equal to the number of linearly independent conditions (No problem with existence!)

The discrete scalar product and norm in analogy with the continuous case is: $\hat{\mathcal{H}}_{N-1} = \hat{\mathcal{H}}_{N-1} + \hat{\mathcal{H}}_{N-1}$

$$\longrightarrow (u, v)_h = \sum_{j=1}^{N-1} q_j < u_j, \hat{H}v_j > h, ||u||^2 = (u, u)_h$$

 $q_j > 0$ and \tilde{H} pos. def. symmetric matrix. Numbering may vary, may include also boundary points.

Def.: Let V_n be space of grid vector functions v satisfying boundary condition $B_n V = 0$. The difference operator Q is <u>semi-bounded</u> if for all $v \in V_n$

$$(v, qv)_n \le \alpha ||v||_n^2$$

Example:

$$\begin{aligned} \frac{d}{dt}u_j &= Du_j, \quad j = 0, 1, \dots, N-1\\ u_N(t) &= 0\\ u_j(0) &= f, \qquad j = 0, \dots, N\\ Du_j &= \begin{cases} D_+u_j & j = 0\\ D_0u_j & j = 1, 2, \dots, N-1 \end{cases} \end{aligned}$$

Redefine scalar product as

$$(u, v)_h = \delta h u_0 v_0 + \sum_{j=1}^{N-1} u_j v_j h = u^T P v.$$

Then

$$(u, Qu) = (u, Du) =$$

$$\delta h u_0 (Du)_0 + \sum_{j=1}^{N-1} u_j (\frac{u_{j+1} - u_{j-1}}{2h}) h = \delta (u_0 u_1 - u_0^2) + u_1 \frac{u_2 - u_0}{2} + u_2 \frac{u_3 - u_1}{2} + \dots$$

$$u_{N-2} (u_{N-1} - u_N - 3) + u_{N-1} (u_N - u_{N-2}) = -\delta u_0^2 + u_0 u_1 (\delta - \frac{1}{2})$$

$$= -\frac{1}{2} u_0 \text{ if } \delta = \frac{1}{2}$$

 \therefore Q semi-bounded (SBP trick!!)

<u>Continuous Problem</u>: Semi-boundedness is shown using integration-by-parts

Semi-discrete Problem: Semi-boundedness is shown by summation-by-parts

Definition:

The problem (2.32) is stable if for F=g=0, $||u||_h \leq ke^{\alpha t}||f||_h$ holds. k and α are constants independent of f and h.

Note that the constants have to be independent of h. The estimate must be independent of grid.

<u>**Theorem**</u>: If Q is semi-bounded, then (2.32) is stable.

Note! No problem with existence and number of boundary conditions and maximal semi-boundedness etc.

Note also!! A non-zero forcing function F is of no problem if ||F|| bounded, see ex (2.29) in book.

Definition The problem 2.32 is strongly stable if it satisfies

$$\begin{split} ||u||^2 &\leq K e^{\alpha t} (||f||^2 + \int_0^t (||F||^2 + ||g||^2) dt))) \\ & (\int_0^T e^{-\alpha \tau} (\frac{1}{\alpha} ||F||^2 + g^2) d\tau) \end{split}$$

K and α are constants independent of $F,\,f,\,g,\,h.$

6 Time-stability

The formal stability definitions allow for exponential growth in time for fixed h, k. (only accurate in the limit)

Boundedness in time is very useful for long-time calculations, we will use the concept <u>time stable</u>.

Definition in the book

The problem (2.32) is time stable if for F = g = 0, there is a unique solution satisfying

$$||u||_h^2 \le k||f||_h$$

k independent of $f,\,h,\,t.$

<u>A better definition</u>

Assume that the PDE for F = g = 0 has the estimate

$$||u|| \le k_c e^{\alpha_c t} ||f||$$

The difference approximation (2.32) is time-stable if it has the estimate

$$||u||_h^2 \le k_d e^{\alpha_d t} ||f||_h$$

where $\alpha_d \leq \alpha_c + \mathcal{O}(h)$.

We will come back to time-stability later once we have developed more theory and skill.



<u>Ex:</u> Time stability (periodic case)

Continuous:

The PDE is an skew-symmetric form

$$u_t = p(u) = (au)_x + a(x)u_x, \ a > 0$$

$$u(1,t) = u(0,t)$$

$$u(x,0) = f(x)$$

$$(u, pa) = (u, (au)_x + au_x) = \int_0^1 u(au)_x + auu_x dx = \underbrace{au^2|_0^1}_{=0 \ periodic} + \underbrace{\int_0^1 - auu_x + auu_x dx}_{=0} \le 0$$

: P Semi bounded $\Rightarrow \frac{d}{dt} ||u||^2 = 0$

<u>Semi discrete</u>

$$(u_j)_t = Qu_j = D_0(a_ju_j) + a_jD_0u_j$$

 $x_j = jh, \ j = 0, ..., N, \ Nh = 1, \ u_j = u_j + N?, \ \text{all } j.$

$$(u, v) = \sum_{j=0}^{N-1} u_j v_j h$$
$$(u, D_0 v)_h = \frac{1}{2} \sum_{j=0}^{N-1} u_j (u_{j+1} - u_{j-1}) =$$
$$= \frac{1}{2} \sum_{j=1}^{N} u_{j-1} v_j - \frac{1}{2} \sum_{j=-1}^{N-2} u_{j+1} v_j$$
$$= \frac{1}{2} \sum_{j=0}^{N-1} u_{j-1} v_j - \frac{1}{2} \sum_{j=0}^{N-1} u_{j+1} v_j = -\frac{1}{2} \sum_{j=0}^{N-1} (u_{j+1} - u_{j-1}) v_j = -(D_0 u, v)$$

 $\therefore D_0$ is skew-symmetric

$$(u, Qu) = \underbrace{(u, D_0(au))}_{=-(D_0u, au)} + (u, aD_0u) = 0$$
$$\frac{d}{dt} ||u||_h^2 = 0$$
$$\frac{du_j}{dt} = Qu_j \text{ time-stable}$$

$$(a, D_0(Au)) + (u, AD_0u) = -(D_0u, Au) + (u, AD_0u) = -(D_0u)^T + Au + u^T A D_0u = 0$$