

## TATA27 Partiella differentialekvationer

### Tentamen 2023-05-30 kl. 8.00–12.00

No aids allowed (except drawing tools, such as rulers, of course). You may write your answers in English or in Swedish, or some mixture thereof.

Each problem is marked *pass* (3 or 2 points) or *fail* (1 or 0 points). For grade  $n \in \{3, 4, 5\}$  you need at least  $n$  passed problems and at least  $3n - 1$  points.

Solutions will be posted on the course webpage afterwards. Good luck!

1. Solve the wave equation  $u_{tt} = c^2 u_{xx}$  for  $0 < x < \pi/2$  and  $t > 0$ , with the boundary conditions  $u_x(0, t) = 0$  and  $u(\pi/2, t) = 0$  and the initial conditions  $u(x, 0) = 2 \cos(3x)$  and  $u_t(x, 0) = 7 \cos(5x)$ .
2. Formulate the weak and strong maximum principles for harmonic functions on a nonempty open set  $\Omega$  in  $\mathbf{R}^n$ .
3. Use the method of characteristics to solve the PDE  $xu_x + (1 + y^2)u_y = u$  with the condition  $u(x, 0) = f(x)$ , where  $f \in C^1(\mathbf{R})$  is some given function.
4. Suppose  $u(x, y, z)$  is harmonic on the open ball  $x^2 + y^2 + z^2 < 4$  and continuous out to the boundary, with the boundary values given by  $u(x, y, z) = \exp(z)$  for  $x^2 + y^2 + z^2 = 4$ . Determine the value  $u(0, 0, 0)$ .
5. Consider the PDE  $u_t + cu_x = \alpha u_{xx}$  where  $\alpha$  and  $c$  are positive constants.
  - (a) Suggest (with motivation) some physical situation which may be modelled by this equation.
  - (b) Find a change of variables of the form  $(\tau, \xi) = (t, f(x, t))$  which reduces the PDE to the heat equation  $u_\tau = \alpha u_{\xi\xi}$ .
6. Consider the following finite element approach to the heat equation  $u_t = u_{xx}$  on the interval  $0 < x < 1$ , with Dirichlet boundary conditions  $u = 0$  at the endpoints: introduce  $N$  nodes  $x_k$  such that  $0 < x_1 < x_2 < \dots < x_N < 1$ , and seek an approximate solution of the form  $u(x, t) = \sum_{k=1}^N c_k(t) \varphi_k(x)$ , where  $\varphi_k(x)$  is the standard “tent-shaped” basis function which is piecewise linear, equals 1 at the node  $x_k$ , and equals 0 at all other nodes.
  - (a) Derive a semi-weak formulation suitable for the FEM approach above. (Multiply by a test function  $\varphi(x)$  and integrate by parts to move a derivative from  $u_{xx}$  to  $\varphi$ .)
  - (b) Show that this gives an ODE system of the form  $A \frac{d\mathbf{c}}{dt} + B\mathbf{c} = \mathbf{0}$  for the vector  $\mathbf{c}(t) = (c_1(t), \dots, c_N(t))^T$ , where  $A$  and  $B$  are symmetric tri-diagonal  $N \times N$  matrices with positive entries on the main diagonal.

## Solutions for TATA27 2023-05-30

1. The basic separated solutions that satisfy the PDE and the boundary conditions are  $u(x, t) = \cos(\omega x) \cos(\omega ct)$  and  $u(x, t) = \cos(\omega x) \sin(\omega ct)$  with  $\omega$  an odd positive integer, and the general solution is a Fourier-type series, a linear combination of these infinitely many basic solutions. However, the particular initial conditions given here are satisfied by a very simple linear combination, where only two terms (with  $\omega = 3$  and  $\omega = 5$ ) are nonzero.

**Answer.**  $u(x, t) = 2 \cos(3x) \cos(3ct) + \frac{7}{5c} \cos(5x) \sin(5ct)$ .

2. The weak maximum principle says that if  $u$  is harmonic on a **bounded** nonempty open set  $\Omega$  and  $u \in C(\overline{\Omega})$ , then the maximum and minimum of  $u$  on  $\overline{\Omega}$  (which exist by the extreme value theorem) are attained on the boundary  $\partial\Omega$ . The strong maximum principle says that if  $\Omega$  moreover is **connected**, then the maximum and minimum are attained *only* on the boundary, not in the interior, unless  $u$  is constant on  $\overline{\Omega}$ .

[A related statement which is also sometimes called the strong maximum principle is that if  $u$  is harmonic on a connected (but not necessarily bounded) open set  $\Omega$  and has a local maximum or minimum at some point in  $\Omega$ , then  $u$  is constant on  $\Omega$ .]

3. For a fixed  $s \in \mathbf{R}$ , the characteristic curve  $(x(t), y(t))$  through the point  $(s, 0)$  is given by  $\dot{x} = x$ ,  $x(0) = s$  and  $\dot{y} = 1 + y^2$ ,  $y(0) = 0$ , hence  $x(t) = se^t$  and  $y(t) = \tan t$ ,  $|t| < \pi/2$ . Along that curve,  $z(t) = u(x(t), y(t))$  satisfies  $\dot{z} = z$  with  $z(0) = f(s)$ , so that  $z(t) = f(s)e^t$ . With  $t = \arctan y$  and  $s = x/e^t = xe^{-\arctan y}$ , this gives the solution  $u = f(s)e^t = f(xe^{-\arctan y})e^{\arctan y}$  (defined in the whole plane).

**Answer.**  $u(x, y) = f(xe^{-\arctan y})e^{\arctan y}$ ,  $(x, y) \in \mathbf{R}^2$ .

4. The mean value property for harmonic functions says that  $u(0, 0, 0)$  equals the average of the values on the boundary sphere  $x^2 + y^2 + z^2 = 4$ . To compute this average, we parametrize the sphere with spherical coordinates

$$(x, y, z) = (r \cos \varphi \sin \theta, r \sin \varphi \sin \theta, r \cos \theta)$$

with  $r = 2$ , which gives  $dS = r \cdot r \sin \theta \cdot d\theta d\varphi = 4 \sin \theta d\theta d\varphi$  (using the scale factors  $r$  and  $r \sin \theta$  for the  $\theta$  and  $\varphi$  directions). This gives the mean value integral

$$\begin{aligned} u(0, 0, 0) &= \frac{\int_S u dS}{\int_S dS} = \frac{\int_{\theta=0}^{\pi} \left( \int_{\varphi=0}^{2\pi} \exp(2 \cos \theta) 4 \sin \theta d\varphi \right) d\theta}{\int_{\theta=0}^{\pi} \left( \int_{\varphi=0}^{2\pi} 4 \sin \theta d\varphi \right) d\theta} \\ &= \frac{8\pi \int_{\theta=0}^{\pi} \exp(2 \cos \theta) \sin \theta d\theta}{8\pi \int_{\theta=0}^{\pi} \sin \theta d\theta} = \frac{\left[ -\frac{1}{2} \exp(2 \cos \theta) \right]_{\theta=0}^{\pi}}{\left[ -\cos \theta \right]_{\theta=0}^{\pi}} \\ &= \frac{-\frac{1}{2}(\exp(-2) - \exp(2))}{2} = \frac{1}{2} \sinh(2). \end{aligned}$$

5. (a) This is a diffusion–advection equation, describing for example diffusion of some chemical (of concentration  $u$ ) in a fluid flowing with constant speed  $c$  in a long tube along the  $x$ -axis. It’s a conservation law  $u_t + J_x = 0$ , where the flux term  $J = -\alpha u_x + cu$  combines Fick’s law of diffusion  $J_1 = -\alpha u_x$  with an advection term  $J_2 = cu$ .
- (b) Inspired by part (a), we “go with the flow” and let  $\xi = x - ct$  (together with  $\tau = t$ ). A short computation with the chain rule verifies that we do get the heat (or diffusion) equation in these new variables.
6. (a) With a test function  $\varphi(x)$  which is zero at the endpoints  $x = 0$  and  $x = 1$ , we get

$$\begin{aligned} 0 &= \int_0^1 (u_t - u_{xx}) \varphi \, dx = \int_0^1 u_t \varphi \, dx - \int_0^1 u_{xx} \varphi \, dx \\ &= \int_0^1 u_t \varphi \, dx - \left( \underbrace{[u_x \varphi]_0^1}_{=0} - \int_0^1 u_x \varphi_x \, dx \right) \\ &= \int_0^1 (u_t \varphi + u_x \varphi_x) \, dx, \end{aligned}$$

so a suitable semi-weak formulation is to require this last integral to be zero for all test functions  $\varphi(x)$  such that  $\varphi(0) = \varphi(1) = 0$ .

- (b) The FEM approximation to the solution is obtained by taking  $u(x, t) = \sum_{k=1}^N c_k(t) \varphi_k(x)$  and requiring the integral above to be zero when  $\varphi = \varphi_m$  (for  $1 \leq m \leq N$ ). Writing dot for  $\frac{d}{dt}$  and prime for  $\frac{d}{dx}$  we get

$$\begin{aligned} 0 &= \int_0^1 \left( \left( \sum_{k=1}^N \dot{c}_k(t) \varphi_k(x) \right) \varphi_m(x) + \left( \sum_{k=1}^N c_k(t) \varphi'_k(x) \right) \varphi'_m(x) \right) dx \\ &= \sum_{k=1}^N \left( \underbrace{\left( \int_0^1 \varphi_k(x) \varphi_m(x) \, dx \right)}_{=A_{mk}} \dot{c}_k(t) + \underbrace{\left( \int_0^1 \varphi'_k(x) \varphi'_m(x) \, dx \right)}_{=B_{mk}} c_k(t) \right) \\ &= \sum_{k=1}^N (A_{mk} \dot{c}_k(t) + B_{mk} c_k(t)). \end{aligned}$$

This sum is entry number  $m$  in the column vector  $A \frac{d\mathbf{c}}{dt} + B\mathbf{c}$ , so the whole vector must be equal to the zero vector, as was to be shown. It’s obvious from the definitions that the matrices  $A$  and  $B$  satisfy  $A_{mk} = A_{km}$  and  $B_{mk} = B_{km}$ , and they are tridiagonal since if the indices  $k$  and  $m$  are more than one step apart, then  $\varphi_k$  and  $\varphi_m$  have disjoint supports, so that  $\varphi_k(x) \varphi_m(x) = 0$  and  $\varphi'_k(x) \varphi'_m(x) = 0$  for all  $x \in [0, 1]$ . Also, it’s clear that  $A_{kk} = \int_0^1 \varphi_k(x)^2 \, dx > 0$  and  $B_{kk} = \int_0^1 \varphi'_k(x)^2 \, dx > 0$  for  $1 \leq k \leq N$ .