## TATA27 Partiella differentialekvationer

## Tentamen 2023-08-14 kl. 14.00-18.00

No aids allowed (except drawing tools, such as rulers, of course). You may write your answers in English or in Swedish, or some mixture thereof.
Each problem is marked pass ( 3 or 2 points) or fail (1 or 0 points). For grade $n \in\{3,4,5\}$ you need at least $n$ passed problems and at least $3 n-1$ points.
Solutions will be posted on the course webpage afterwards. Good luck!

1. Solve the heat equation $u_{t}=u_{x x}$ for $0<x<\pi$ and $t>0$, with the initial condition $u(x, 0)=x(\pi-x)$ and the Dirichlet boundary conditions $u(0, t)=$ $u(\pi, t)=0$. (Hint: The solution takes the form of a Fourier-type series.)
2. Determine all functions $f(r), r>0$, such that $u(x, y)=f\left(\sqrt{x^{2}+y^{2}}\right)$ is harmonic on $\mathbf{R}^{2} \backslash\{(0,0)\}$.
3. Consider the wave equation $u_{t t}=\frac{1}{4} u_{x x}$ for $x \in \mathbf{R}$ and $t>0$, with initial conditions

$$
u(x, 0)=\left\{\begin{array}{ll}
\sin x, & \text { for } 0<x<\pi, \\
0, & \text { otherwise },
\end{array} \quad u_{t}(x, 0)=0 \quad \text { for } x \in \mathbf{R} .\right.
$$

Let $u(x, t)$ be the (weak) solution given by d'Alembert's formula. Draw the graph of the function $x \mapsto u(x, 1)$. (In other words, draw the shape of the wave after one unit of time.)
4. Use the method of characteristics to find $u(x, y)$ such that $x^{2} u_{x}+u_{y}=u$ and $u(x, 0)=\sin x$. In what region of the plane is the problem's solution $u$ determined?
5. Poisson's integral formula for the unit ball in $\mathbf{R}^{n}$ says that if $u$ is harmonic on the open ball $B(\mathbf{0}, 1)$ and continuous on the closed ball $\overline{B(\mathbf{0}, 1)}$, then

$$
u(\mathbf{a})=f_{S(0,1)} \frac{1-|\mathbf{a}|^{2}}{|\mathbf{x}-\mathbf{a}|^{n}} u(\mathbf{x}) d S(\mathbf{x}) \quad \text { for } \mathbf{a}=B(\mathbf{0}, 1)
$$

where $S(\mathbf{0}, 1)$ is the unit sphere, and the symbol $f$ denotes the mean value integral. Use this to derive Poisson's integral formula for a ball $B(\mathbf{0}, r)$ in $\mathbf{R}^{n}$ of arbitrary radius $r>0$.
6. Formulate the weak maximum principle for the heat equation $u_{t}=\Delta u$, where $t \in(0, T)$ and $\mathbf{x} \in \Omega$ (a bounded domain in $\mathbf{R}^{n}$ ). Use this to show uniqueness of solutions to the initial-boundary value problem where the initial temperature $u(\mathbf{x}, 0)$ and the boundary values $u(\mathbf{x}, t), \mathbf{x} \in \partial \Omega$, are prescribed.

## Solutions for TATA27 2023-08-14

1. The basic separated solutions satisfying the PDE and the boundary conditions are $u_{n}(x, t)=\sin (n x) e^{-n^{2} t}$ for integers $n \geq 1$, and the sought solution takes the form $u(x, t)=\sum_{n=1}^{\infty} c_{n} u_{n}(x, t)$. In order to satisfy the initial condition, we need $x(\pi-x)=u(x, 0)=\sum_{n=1}^{\infty} c_{n} u_{n}(x, 0)=\sum_{n=1}^{\infty} c_{n} \sin (n x)$ for $0<x<\pi$. We multiply this by $\sin (k x)$ and integrate from 0 to $\pi$, using that $\int_{0}^{\pi} \sin (n x) \sin (k x) d x=0$ for $k \neq n$, to obtain

$$
c_{k} \underbrace{\int_{0}^{\pi} \sin ^{2}(k x) d x}_{=\pi / 2}=\int_{0}^{\pi} x(\pi-x) \sin (k x) d x
$$

so that

$$
\begin{aligned}
c_{k} & =\frac{2}{\pi} \int_{0}^{\pi} x(\pi-x) \sin (k x) d x \\
& =\frac{2}{\pi}\left[\left(\pi x-x^{2}\right) \frac{-\cos (k x)}{k}-(\pi-2 x) \frac{-\sin (k x)}{k^{2}}+(-2) \frac{\cos (k x)}{k^{3}}\right]_{0}^{\pi} \\
& =\frac{4\left(1-(-1)^{k}\right)}{\pi k^{3}}
\end{aligned}
$$

## Answer.

$$
u(x, t)=\sum_{n=1}^{\infty} \frac{4\left(1-(-1)^{n}\right)}{\pi n^{3}} \sin (n x) e^{-n^{2} t}
$$

or equivalently, since only odd-numbered $c_{n}=c_{2 m+1}$ are nonzero,

$$
u(x, t)=\sum_{m=0}^{\infty} \frac{8}{\pi(2 m+1)^{3}} \sin ((2 m+1) x) e^{-(2 m+1)^{2} t}
$$

2. If you remember the formula for the Laplacian in polar coordinates, $\Delta u=$ $u_{r r}+\frac{1}{r^{2}} u_{\varphi \varphi}+\frac{1}{r} u_{r}$, you can use that to obtain $f^{\prime \prime}(r)+\frac{1}{r} f^{\prime}(r)=0$. Otherwise, just compute $u_{x}(x, y)=f^{\prime}\left(\sqrt{x^{2}+y^{2}}\right) \cdot x / \sqrt{x^{2}+y^{2}}$, and so on, to derive that same ODE. Multiplication by the integrating factor $r$ gives $\left(r f^{\prime}(r)\right)^{\prime}=0$, and after two integrations we find $f(r)=A \ln r+B$ with arbitrary constants $A$ and $B$.

Answer. $f(r)=A \ln r+B, r>0$.
3. Since $c=\sqrt{1 / 4}=1 / 2$, d'Alembert's formula says that

$$
u(x, t)=\frac{\varphi\left(x-\frac{1}{2} t\right)+\varphi\left(x+\frac{1}{2} t\right)}{2}
$$

where $\varphi(x)=u(x, 0)$, and in particular

$$
u(x, 1)=\frac{\varphi\left(x-\frac{1}{2}\right)+\varphi\left(x+\frac{1}{2}\right)}{2} .
$$

To draw the graph, we can first plot the two terms $\frac{1}{2} \varphi\left(x \pm \frac{1}{2}\right)$ separately:


Then adding them up (graphically) gives the answer:

4. For a fixed $s \in \mathbf{R}$, the characteristic curve $(x(t), y(t))$ through the point $(s, 0)$ is given by $\dot{x}=x^{2}$ and $\dot{y}=1$, with initial conditions $x(0)=s$ and $y(0)=0$. Hence, $x(t)=s /(1-s t)$ and $y(t)=t$. (The solution for $x(t)$ is valid for all $t \in \mathbf{R}$ if $s=0$, but only for $t<1 / s$ if $s>0$ and for $t>1 / s$ if $s<0$.) Along such a curve, $z(t)=u(x(t), y(t))$ satisfies $\dot{z}=z$ with $z(0)=\sin (s)$, so that $z(t)=\sin (s) e^{t}$. With $t=y$ and $s=x /(1+x t)=x /(1+x y)$, this gives $u=\sin (s) e^{t}=\sin (x /(1+x y)) e^{y}$. The characteristics computed above, emanating from points on the $x$-axis, fill out the region $x y>-1$ (between the two dotted curves in the figure below), so this is the region where the solution $u$ is determined by the given conditions. If we approach either of the dotted curves from inside this region, for example along a line $x=$ const., the solution is seen to oscillate wildly, so it cannot be continued up to (or across) the boundary in a nice way.
Answer. $u(x, y)=\sin \left(\frac{x}{1+x y}\right) e^{y}$, for $x y>-1$.

5. Let $A_{n}$ be the ( $n-1$ )-dimensional surface area of the unit sphere $S(\mathbf{0}, 1)$ in $R^{n}$. Suppose $u$ is harmonic on $B(\mathbf{0}, r)$ and continuous out to the boundary. Then $v(\mathbf{x})=u(r \mathbf{x})$ is harmonic on the unit ball $B(\mathbf{0}, 1)$ and continuous out to the boundary. (Rescaling a harmonic function like this doesn't destroy harmonicity; this can be shown by calculation using the chain rule, or argued more abstractly, for example by noticing that the rescaled function still has the mean value property.) Thus, for $\mathbf{a} \in B(\mathbf{0}, r)$, we can apply the formula in the problem to the rescaled function $v$, with $\mathbf{b}=\frac{1}{r} \mathbf{a} \in B(\mathbf{0}, 1)$ playing the role of $\mathbf{a}$, and writing $\mathbf{y}$ instead of $\mathbf{x}$, to obtain the value

$$
\begin{aligned}
u(\mathbf{a}) & =v\left(\frac{1}{r} \mathbf{a}\right)=v(\mathbf{b})=f_{S(0,1)} \frac{1-|\mathbf{b}|^{2}}{|\mathbf{y}-\mathbf{b}|^{n}} v(\mathbf{y}) d S(\mathbf{y}) \\
& =f_{S(\mathbf{0}, 1)} \frac{1-\left|\frac{1}{r} \mathbf{a}\right|^{2}}{\left|\mathbf{y}-\frac{1}{r} \mathbf{a}\right|^{n}} v(\mathbf{y}) d S(\mathbf{y})=\frac{1}{A_{n}} \int_{S(\mathbf{0}, 1)} \frac{\frac{1}{r^{2}}}{\frac{1}{r^{n}}} \frac{r^{2}-|\mathbf{a}|^{2}}{|r \mathbf{y}-\mathbf{a}|^{n}} u(r \mathbf{y}) d S(\mathbf{y}) .
\end{aligned}
$$

Under the change of variables $\mathbf{x}=r \mathbf{y}$, the surface area element scales as $d S(\mathbf{x})=r^{n-1} d S(\mathbf{y})$, and the new region of integration is the sphere $S(\mathbf{0}, r)$, whose surface area is $A_{n} r^{n-1}$, so we get

$$
\begin{aligned}
u(\mathbf{a}) & =\frac{1}{A_{n}} \int_{S(\mathbf{0}, r)} \frac{r^{n}}{r^{2}} \frac{r^{2}-|\mathbf{a}|^{2}}{|\mathbf{x}-\mathbf{a}|^{n}} u(\mathbf{x}) \frac{d S(\mathbf{x})}{r^{n-1}} \\
& =\frac{1}{A_{n} r^{n-1}} \int_{S(\mathbf{0}, r)} \frac{r^{n-2}\left(r^{2}-|\mathbf{a}|^{2}\right)}{|\mathbf{x}-\mathbf{a}|^{n}} u(\mathbf{x}) d S(\mathbf{x}) \\
& =f_{S(\mathbf{0}, r)} \frac{r^{n-2}\left(r^{2}-|\mathbf{a}|^{2}\right)}{|\mathbf{x}-\mathbf{a}|^{n}} u(\mathbf{x}) d S(\mathbf{x}),
\end{aligned}
$$

which is the sought formula.
6. Suppose $u(\mathbf{x}, t)$ solves the heat equation on the set $\Omega \times(0, T)$ and is continuous out to the boundary. Then the weak maximum principle says that the maximum value of $u$ on the compact set $\bar{\Omega} \times[0, T]$ is attained at a point on the parabolic boundary of $\bar{\Omega} \times[0, T]$, i.e., at a point in that set such that either $t=0$ or $\mathbf{x} \in \partial \Omega$. (And applying this to the function $-u$ shows that the minimum of $u$ is likewise attained on the parabolic boundary.)
Now if $u_{1}$ and $u_{2}$ are two solutions to the same initial-boundary value problem, then their difference $u=u_{1}-u_{2}$ is a solution to the heat equation which is zero on the parabolic boundary, and hence by the weak maximum (and minimum) principle, it must be zero throughout $\bar{\Omega} \times[0, T]$. So $u_{1}=u_{2}$, which means that the initial-boundary value problem can have at most one solution.

