

Transform theory 2022-08-18 – Solutions

1. (a) No. The function does not tend to zero as $|\omega| \rightarrow \infty$ so the Riemann-Lebesgue lemma is not satisfied.
- (b) No. For example $f(x) = \frac{1}{x}$, $x > 1$, and $f(x) = 0$ elsewhere, is a function such that $f \in L^2(\mathbf{R})$ but $f \notin L^1(\mathbf{R})$.
- (c) Yes. A continuous function is bounded on a closed interval, so let K be such that $|u(x)| \leq K$ on $[-\pi, \pi]$. Then

$$\int_{-\pi}^{\pi} |u(x)| dx \leq K \int_{-\pi}^{\pi} dx = 2\pi K < \infty$$

and

$$\int_{-\pi}^{\pi} |u(x)|^2 dx \leq K^2 \int_{-\pi}^{\pi} dx = 2\pi K^2 < \infty,$$

so $u \in L^1(-\pi, \pi)$ and $u \in L^2(-\pi, \pi)$.

- (d) Considering that $u(x) = (1+x^2)^{-1}$ is a C^1 -funktion on $] -\pi, \pi[$, continuous on $[-\pi, \pi]$ with $u(-\pi) = u(\pi)$, it follows from a well known theorem that the convergence of the Fourier series is uniform.
- (e) No. For instance, $u[k] = 1$ has the Z transform $z/(z-1)$ for $|z| > 1$ and clearly this expression is unbounded as $z \rightarrow 1$.

Answer: No, No, Yes, Yes, No.

2. We take the Z transform of the equation and find that

$$z^2 U(z) - z^2 u[0] - zu[1] + zU(z) - zu[0] - 2U(z) = \frac{z}{z-1},$$

where we assume that (at least) $|z| > 1$. Reformulating this equation, we find that

$$\begin{aligned} U(z)(z^2 + z - 2) &= 3z^2 + \frac{10z}{3} + \frac{z}{z-1} \quad \Leftrightarrow \quad \frac{U(z)}{z} = \frac{3z + 10/3}{(z-1)(z+2)} + \frac{1}{(z-1)^2(z+2)} \\ \Leftrightarrow \quad \frac{U(z)}{z} &= \frac{19/9}{z-1} + \frac{8/9}{z+2} + \frac{1/3}{(z-1)^2} - \frac{1/9}{z-1} + \frac{1/9}{z+2} = \frac{1/3}{(z-1)^2} + \frac{1}{z+2} + \frac{2}{z-1}, \end{aligned}$$

where we assume that $|z| > 2$. From a table we find that

$$\mathcal{Z}(k) = \frac{z}{(z-1)^2}, \quad \mathcal{Z}((-2)^k) = \frac{z}{z+2}, \quad \text{and} \quad \mathcal{Z}(1) = \frac{z}{z-1},$$

so linearity and uniqueness imply that

$$u[k] = \frac{k}{3} + (-2)^k + 2.$$

Answer: $u[k] = \frac{k}{3} + (-2)^k + 2$, $k = 0, 1, 2, 3, \dots$

3. Clearly $u \in E$. We find that, for $k \neq 0$,

$$\begin{aligned} c_k &= \frac{1}{2\pi} \int_{-\pi}^{\pi} u(x) e^{-ikx} dx = \frac{1}{2\pi} \int_0^{\pi} (\pi - x) e^{-ikx} dx \\ &= \frac{1}{2\pi} \left(\left[(\pi - x) \frac{e^{-ikx}}{-ik} \right]_0^{\pi} - \frac{1}{ik} \int_0^{\pi} e^{-ikx} dx \right) \\ &= \frac{1}{2\pi} \left(\frac{\pi}{ik} - \frac{1}{k^2} [e^{-ikx}]_0^{\pi} \right) = \frac{-i}{2k} + \frac{1 - (-1)^k}{2\pi k^2} = \frac{1 - ik\pi - (-1)^k}{2\pi k^2} \end{aligned}$$

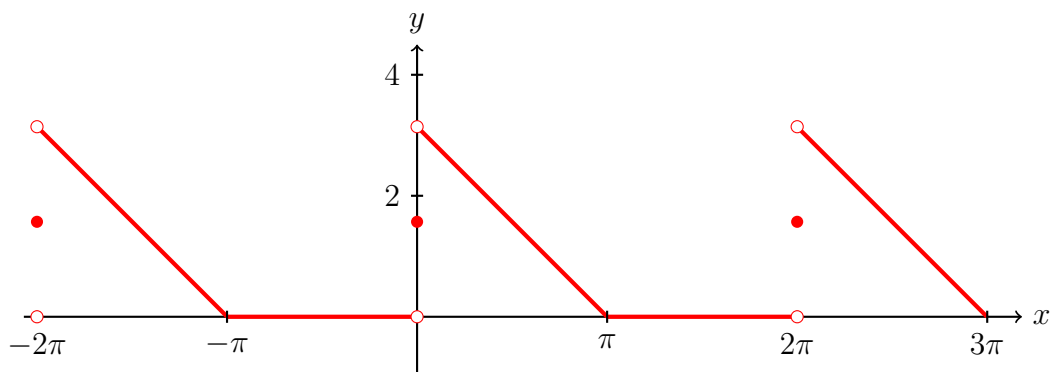
and

$$c_0 = \frac{1}{2\pi} \int_0^{\pi} (\pi - x) dx = \frac{\pi}{4}.$$

Hence

$$u(x) \sim \frac{\pi}{4} + \sum_{k \neq 0} \left(\frac{1 - ik\pi - (-1)^k}{2\pi k^2} \right) e^{ikx}.$$

Furthermore, the function u is differentiable for $x \neq m\pi$, $m \in \mathbf{Z}$, and the function has right- and lefthand derivatives at $x = m\pi$ (it is piecewise straight lines). Hence – by Dirichlet's theorem – the Fourier series is convergent and converges to $u(x)$ for every $x \neq m\pi$, to 0 at $x = (2l + 1)\pi$ ($l \in \mathbf{Z}$) and to $\pi/2$ for $x = 2l\pi$ ($l \in \mathbf{Z}$).



For $k \neq 0$, note that

$$c_k = \frac{1 - ik\pi - (-1)^k}{2\pi k^2} = \frac{1 - (-1)^k}{2\pi k^2} - i \frac{1}{2k}.$$

With $c_0 = \pi/4$, we therefore have

$$|c_k|^2 = \frac{(1 - (-1)^k)^2}{4\pi^2 k^4} + \frac{1}{4k^2}, \quad k \neq 0, \quad |c_0|^2 = \frac{\pi^2}{16},$$

so

$$\sum_{k \in \mathbf{Z}} |c_k|^2 = \frac{\pi^2}{16} + 2 \sum_{k=1}^{\infty} \frac{(1 - (-1)^k)^2}{4\pi^2 k^4} + \frac{1}{4k^2} = \frac{\pi^2}{16} + 2 \sum_{k=1}^{\infty} \frac{(1 - (-1)^k)^2 + \pi^2 k^2}{4\pi^2 k^4}$$

since $|c_{-k}|^2 = |c_k|^2$. By Parseval's identity,

$$\sum_{k \in \mathbf{Z}} |c_k|^2 = \frac{1}{2\pi} \int_{-\pi}^{\pi} |u(x)|^2 dx = \frac{1}{2\pi} \int_0^{\pi} (\pi - x)^2 dx = \frac{\pi^2}{6}.$$

Rearranging the terms, we obtain

$$2 \sum_{k=1}^{\infty} \frac{(1 - (-1)^k)^2 + \pi^2 k^2}{4\pi^2 k^4} = \frac{\pi^2}{6} - \frac{\pi^2}{16} = \frac{5\pi^2}{48},$$

so

$$\sum_{k=1}^{\infty} \frac{(1 - (-1)^k)^2 + \pi^2 k^2}{k^4} = \frac{5\pi^4}{24},$$

which was what we set out to prove.

Answer: $u(x) \sim \frac{\pi}{4} + \sum_{k \neq 0} \left(\frac{1 - ik\pi - (-1)^k}{2\pi k^2} \right) e^{ikx}$; see above.

4. Assuming that $u, u', u'' \in G$, we take the Fourier transform to find that

$$(i\omega)^2 U(\omega) - i\omega U(\omega) - 6U(\omega) = \mathcal{F}(25e^{-2x}H(x)) = \frac{25}{2 + i\omega}.$$

With $s = i\omega$, we see that $s^2 - s - 6 = (s - 3)(s + 2)$, so

$$U(\omega) = \frac{25}{(2 + i\omega)^2(i\omega - 3)} = -\frac{5}{(2 + i\omega)^2} - \frac{1}{3 - i\omega} - \frac{1}{2 + i\omega}.$$

Notice that

$$\frac{d}{d\omega} \left(\frac{1}{2 + i\omega} \right) = -i \frac{1}{(2 + i\omega)^2}$$

so

$$\frac{1}{(2 + i\omega)^2} = i \frac{d}{d\omega} \left(\frac{1}{(2 + i\omega)} \right) = \mathcal{F}(xe^{-2x}H(x)).$$

From a table, we also find that

$$\mathcal{F}(e^{-2x}H(x)) = \frac{1}{2 + i\omega} \quad \text{and} \quad \mathcal{F}(e^{3x}H(-x)) = \frac{1}{3 - i\omega},$$

so by uniqueness and linearity,

$$\begin{aligned} y(x) &= -5xe^{-x}H(x) - e^{3x}H(-x) - e^{-2x}H(x) = -(5x + 1)e^{-2x}H(x) - e^{3x}H(-x) \\ &= \begin{cases} -e^{-3x}, & x < 0. \\ -(5x + 1)e^{-2x}, & x \geq 0. \end{cases} \end{aligned}$$

This function and its derivatives up to order 2 are absolutely integrable.

Answer: $y(x) = -(5x + 1)e^{-2x}H(x) - e^{3x}H(-x)$.

5. We're looking for exponentially bounded solutions, so we assume that $|u(t)| \leq Ke^{at}$ for some constants $a, C > 0$. Taking the Laplace transform of the equation, we obtain that

$$sU(s) - u(0) + U(s) + \frac{1}{s}U(s) = \int_0^1 e^{-st} dt = \frac{1 - e^{-s}}{s}$$

so

$$U(s) \left(s + 1 + \frac{1}{s} \right) = U(s) \frac{s^2 + s + 1}{s} = 1 + \frac{1 - e^{-s}}{s}.$$

Hence,

$$U(s) = \frac{s}{s^2 + s + 1} + \frac{1 - e^{-s}}{s^2 + s + 1}.$$

Since

$$s^2 + s + 1 = (s + 1/2)^2 + 3/4,$$

we rewrite

$$\frac{s}{s^2 + s + 1} = \frac{s + 1/2}{(s + 1/2)^2 + 3/4} - \frac{1/2}{(s + 1/2)^2 + 3/4}$$

and

$$\frac{1}{s^2 + s + 1} = \frac{2}{\sqrt{3}} \frac{\sqrt{3}/2}{(s + 1/2)^2 + 3/4}.$$

From a table,

$$\begin{aligned} \mathcal{L}(\cos at) &= \frac{s}{s^2 + a^2}, & \mathcal{L}(\sin at) &= \frac{a}{s^2 + a^2}, & \mathcal{L}(e^{-at}f(t)) &= F(s + a), \\ \mathcal{L}(H(t - t_0)f(t - t_0)) &= F(s)e^{-st_0}, \end{aligned}$$

so

$$\begin{aligned} \frac{s + 1/2}{(s + 1/2)^2 + 3/4} - \frac{1/2}{(s + 1/2)^2 + 3/4} &= \mathcal{L}\left(e^{-t/2} \cos \frac{t\sqrt{3}}{2}\right) - \frac{1}{2} \frac{2}{\sqrt{3}} \mathcal{L}\left(e^{-t/2} \sin \frac{t\sqrt{3}}{2}\right) \\ &= \mathcal{L}\left(e^{-t/2} \cos \frac{t\sqrt{3}}{2} - \frac{1}{\sqrt{3}} e^{-t/2} \sin \frac{t\sqrt{3}}{2}\right), \end{aligned}$$

$$\frac{2}{\sqrt{3}} \frac{\sqrt{3}/2}{(s + 1/2)^2 + 3/4} = \mathcal{L}\left(\frac{2}{\sqrt{3}} e^{-t/2} \sin \frac{t\sqrt{3}}{2}\right),$$

and

$$\frac{2}{\sqrt{3}} e^{-s} \frac{\sqrt{3}/2}{(s + 1/2)^2 + 3/4} = \mathcal{L}\left(\frac{2}{\sqrt{3}} e^{-(t-1)/2} \sin \frac{(t-1)\sqrt{3}}{2} H(t-1)\right),$$

Hence

$$\begin{aligned} y(t) &= e^{-t/2} \left(\cos \frac{t\sqrt{3}}{2} - \frac{1}{\sqrt{3}} \sin \frac{t\sqrt{3}}{2} + \frac{2}{\sqrt{3}} \sin \frac{t\sqrt{3}}{2} \right) \\ &\quad - \frac{2}{\sqrt{3}} e^{-(t-1)/2} \sin \left(\frac{(t-1)\sqrt{3}}{2} \right) H(t-1) \end{aligned}$$

by uniqueness and linearity.

Answer: $y(t) = e^{-t/2} \left(\cos \frac{t\sqrt{3}}{2} + \frac{1}{\sqrt{3}} \sin \frac{t\sqrt{3}}{2} \right) - \frac{2}{\sqrt{3}} e^{-(t-1)/2} \sin \left(\frac{(t-1)\sqrt{3}}{2} \right) H(t-1),$
 $t \geq 0.$

6. There are at least two different possibilities for solving this problem.

Alternative 1: consider that

$$\frac{\sin^2 2\omega}{1 + \omega^2} = \frac{1 - \cos 4\omega}{2(1 + \omega^2)} = \frac{1}{2} \cdot \frac{1 - \cos 4\omega}{i\omega} \cdot \frac{i\omega}{1 + \omega^2} = \frac{1}{2} \cdot F(\omega) \cdot \overline{G(\omega)},$$

where

$$F(\omega) = \frac{1 - \cos 4\omega}{i\omega} = \mathcal{F} \left(\underbrace{\frac{1}{2} \operatorname{sgn}(x) (H(x+4) - H(x-4))}_{=f(x)} \right)$$

and

$$G(\omega) = \frac{-i\omega}{1 + \omega^2} = \mathcal{F} \left(\underbrace{\frac{1}{2} \operatorname{sgn}(x) e^{-|x|}}_{=g(x)} \right).$$

We note that f and g both belong to $L^1(\mathbf{R}) \cap L^2(\mathbf{R})$, so by Plancherel's identity

$$\begin{aligned} \int_{-\infty}^{\infty} F(\omega) \overline{G(\omega)} d\omega &= 2\pi \int_{-\infty}^{\infty} f(x) \overline{g(x)} dx = \frac{\pi}{4} \int_{-4}^4 \operatorname{sgn}(x)^2 e^{-|x|} dx = \frac{\pi}{2} \int_0^4 e^{-x} dx \\ &= \frac{\pi}{2} (1 - e^{-4}). \end{aligned}$$

Alternative 2: using the same trigonometric identity as in the previous alternative, we find that

$$\int_{-\infty}^{\infty} \frac{\sin^2 2\omega}{1 + \omega^2} d\omega = \frac{1}{2} \int_{-\infty}^{\infty} \frac{d\omega}{1 + \omega^2} - \frac{1}{2} \int_{-\infty}^{\infty} \frac{\cos 4\omega}{1 + \omega^2} d\omega.$$

The first integral in the right-hand side is well-known:

$$\int_{-\infty}^{\infty} \frac{d\omega}{1 + \omega^2} = [\arctan \omega]_{-\infty}^{\infty} = \pi.$$

In the second integral, we let $\xi = 4\omega$ so that

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{\cos 4\omega}{1 + \omega^2} d\omega &= \frac{1}{4} \int_{-\infty}^{\infty} \frac{\cos \xi}{1 + (\xi/4)^2} d\xi = \frac{1}{4} \int_{-\infty}^{\infty} \frac{16 \cos \xi}{16 + \xi^2} d\xi = 4 \int_{-\infty}^{\infty} \frac{\cos \xi}{16 + \xi^2} d\xi \\ &= 4 \int_{-\infty}^{\infty} \frac{\cos \xi}{16 + \xi^2} d\xi = 4 \int_{-\infty}^{\infty} \frac{\cos \xi}{16 + \xi^2} d\xi \\ &= 4 \operatorname{Re} \left(\int_{-\infty}^{\infty} \frac{e^{-i\xi \cdot 1}}{16 + \xi^2} d\xi \right) = 4 \operatorname{Re} \left(\mathcal{F} \left(\frac{1}{4^2 + \xi^2} \right) (x = 1) \right) \\ &= 4 \operatorname{Re} \left(\frac{\pi}{4} e^{-4 \cdot |1|} \right) = \pi e^{-4}, \end{aligned}$$

where we used the fact that $\frac{\cos \xi}{16 + \xi^2}$ is real valued and that $\operatorname{Re}(e^{-i\xi}) = \cos \xi$. Hence

$$\int_{-\infty}^{\infty} \frac{\sin^2 2\omega}{1 + \omega^2} d\omega = \frac{1}{2} (1 - e^{-4}).$$

Answer: $\frac{1}{2} (1 - e^{-4})$.

7. Note that $e^{-x^2/2}$ and $xe^{-x^2/2}$ belong to $G(\mathbf{R})$, so if $u(x) = e^{-x^2/2}$ and $U(\omega) = \mathcal{F}u(\omega)$, then

$$\begin{aligned} U'(\omega) &= -i \mathcal{F}(xe^{-x^2/2})(\omega) = -i \int_{-\infty}^{\infty} xe^{-x^2/2} e^{-i\omega x} dx \\ &= / \text{I.B.P.} / = -i \left[-e^{-x^2} e^{-i\omega x} \right]_{-\infty}^{\infty} - i \int_{-\infty}^{\infty} \frac{-i\omega}{2} e^{-x^2/2} e^{-i\omega x} dx = -\omega U(\omega). \end{aligned}$$

So U must satisfy

$$U'(\omega) + \omega U(\omega) = 0 \quad \Leftrightarrow \quad \frac{d}{d\omega} \left(e^{\omega^2/2} U(\omega) \right) = 0 \quad \Leftrightarrow \quad U(\omega) = C e^{-\omega^2/2}.$$

However, we can only have one Fourier transform so we need to find a value for C . It is clear that

$$U(0) = \int_{-\infty}^{\infty} e^{-x^2/2} dx = / x = t\sqrt{2} / = \sqrt{2} \int_{-\infty}^{\infty} e^{-t^2} dt = \sqrt{2}\sqrt{\pi} = \sqrt{2\pi}.$$

Therefore $C = U(0) = \sqrt{2\pi}$ and we have shown that

$$\mathcal{F} \left(e^{-x^2/2} \right) (\omega) = \sqrt{2\pi} e^{-\omega^2/2}.$$

Answer: see above.